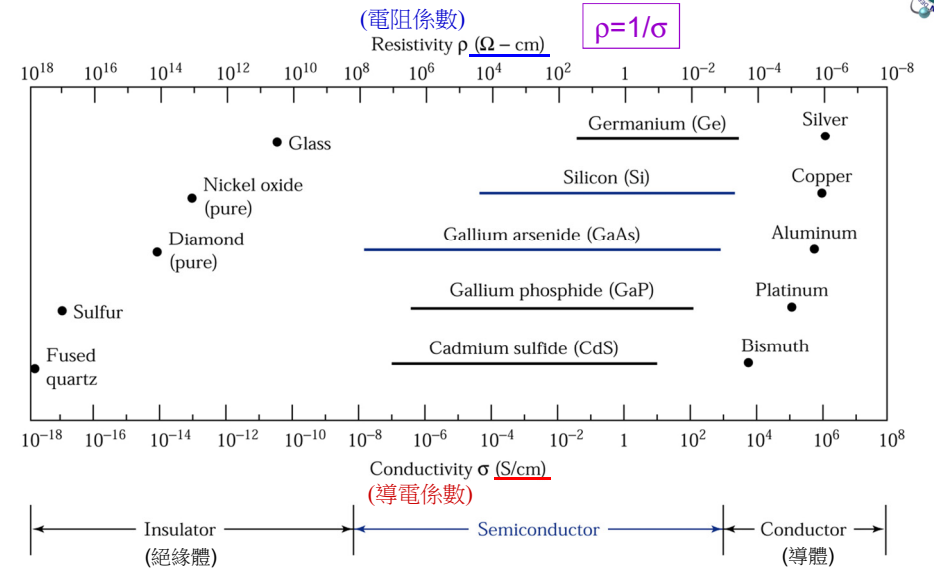




Semiconductor Basics

(半導體基礎)



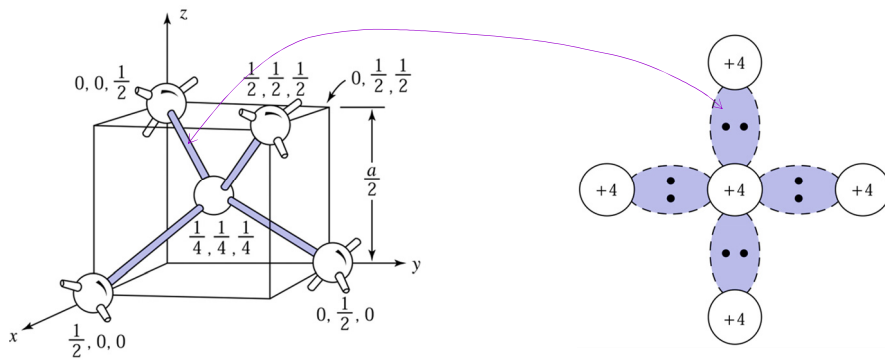
1, jswu

2, jswu

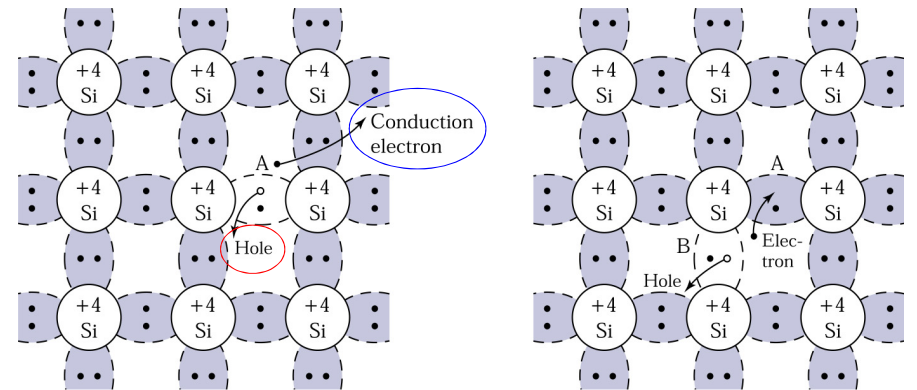
Conduction Electron (導電電子) and Hole (電洞)



- The electrons shared for atomic bonding in semiconductors can not freely move!



- Intrinsic** (本質的) **semiconductor**: a pure semiconductor without *doping* (摻雜).
 \therefore Intrinsic semiconductor \equiv Undoped (未摻雜的) semiconductor.



$T > 0\text{ K}$ (絕對零度)

Electron (電子): a particle of a negative charge ($-1.6 \times 10^{-19}\text{ C}$).

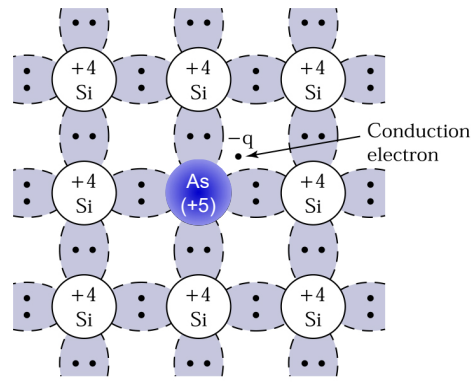
Hole (電洞): a fictitious (虛構的) particle of a positive charge ($+1.6 \times 10^{-19}\text{ C}$).

3, jswu

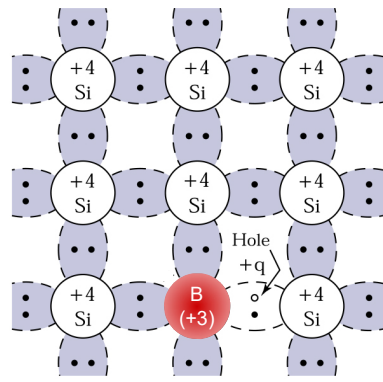
4, jswu



- ◆ **Extrinsic** (外質的) **semiconductor**: a semiconductor doped with impurity atoms (雜質原子)
 - ∴ Extrinsic semiconductor ≡ Doped (摻雜的) semiconductor.
- ◆ The impurity atoms are called **dopants** (摻雜物) when they are used for doping.
- ◆ 2 types of dopants:
 - **donor** (施體) for n-type semiconductors, and
 - **acceptor** (受體) for p type semiconductors.



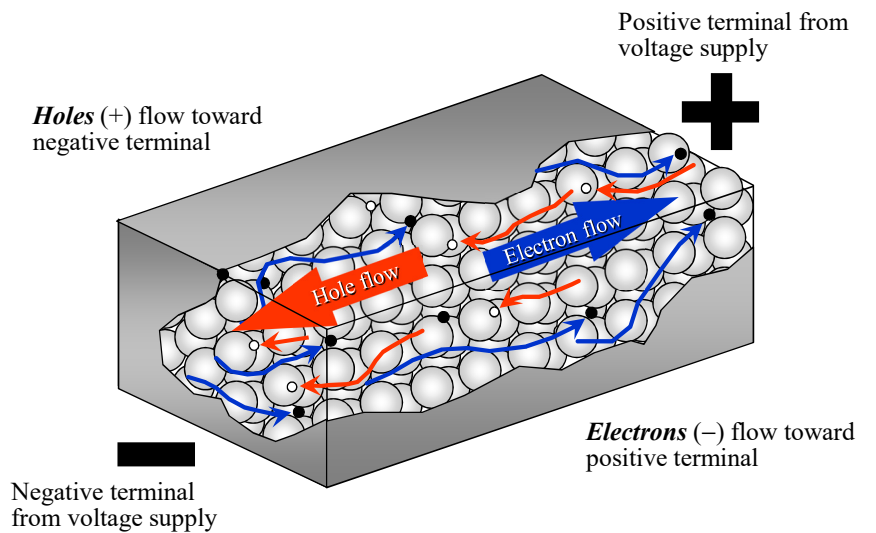
n-type Si with donor (As)



p-type Si with acceptor (B)



Flow of Carriers (載子) in Semiconductor



Holes (+) flow toward negative terminal

Positive terminal from voltage supply

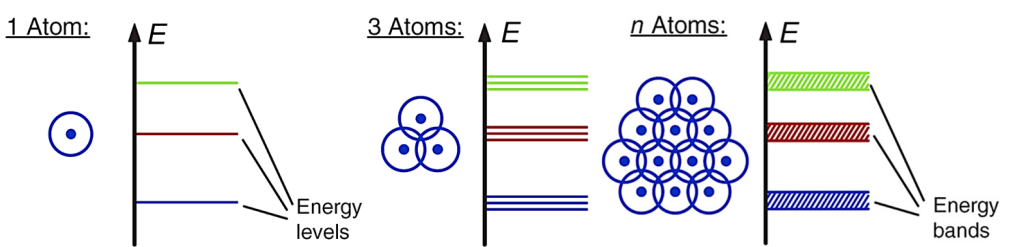
Electrons (-) flow toward positive terminal

Negative terminal from voltage supply



Energy Band (能帶)

Formation of energy bands in solid

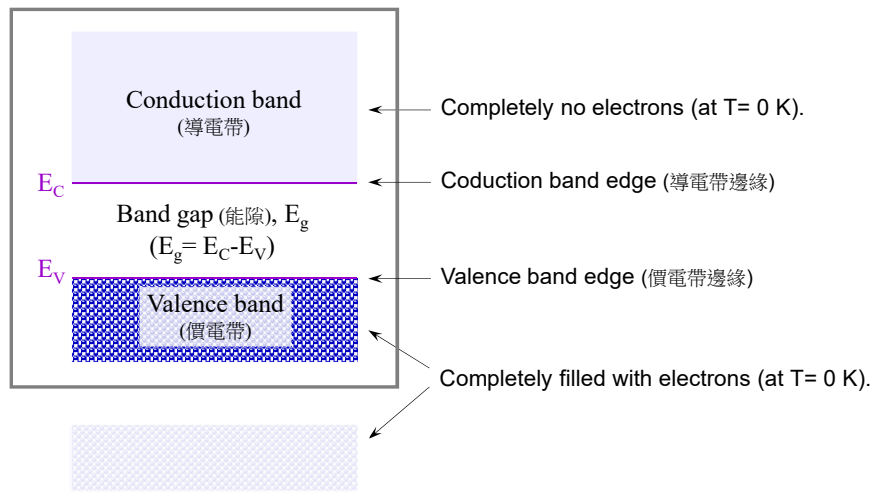


For a solid, a large amount of closely-located energy levels form an energy band.



Energy band diagram (能帶圖) of a semiconductor

Energy band diagram (能帶圖) → E-x relation



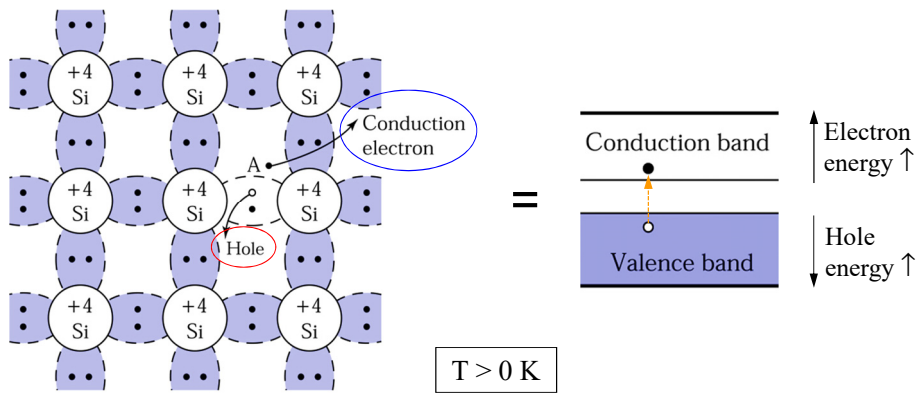
← Completely no electrons (at T= 0 K).

← Coduction band edge (導電帶邊緣)

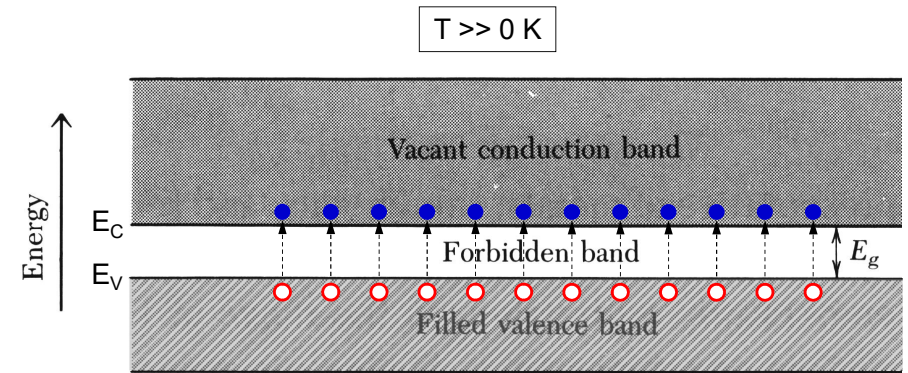
← Valence band edge (價電帶邊緣)

← Completely filled with electrons (at T= 0 K).

- ◆ Electron (電子): a particle of a negative charge (-1.6×10^{-19} C).
Hole (電洞): a fictitious (虛構的) particle of a positive charge ($+1.6 \times 10^{-19}$ C).
- ◆ Free electrons are in the conduction band and holes are in the valence band.



- ◆ As $T \uparrow \uparrow$, more electrons are *thermally excited* (熱激發) from the valence band to the conduction band.
- ◆ Electron concentration (電子濃度) n = Hole concentration (電洞濃度) p .



Band gap & electron occupancy (電子佔居) at 0 K



Insulator
(絕緣體)

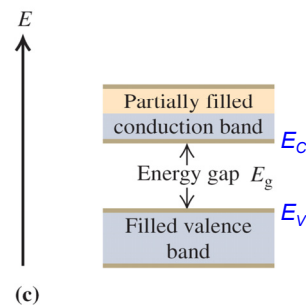
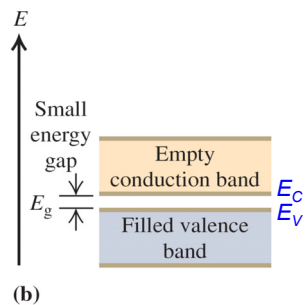
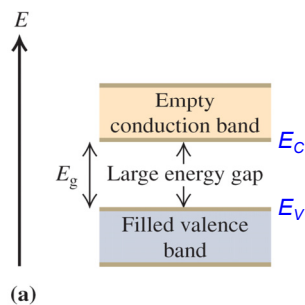
Semiconductor

Conductor
(導體)

In an insulator at absolute zero, there are no electrons in the conduction band.

A semiconductor has the same band structure as an insulator but a smaller gap between the valence and conduction bands.

A conductor has a partially filled conduction band.

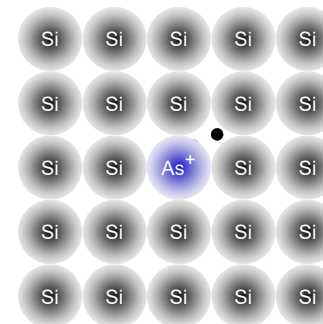


Extrinsic semiconductor: when a donor atom is present



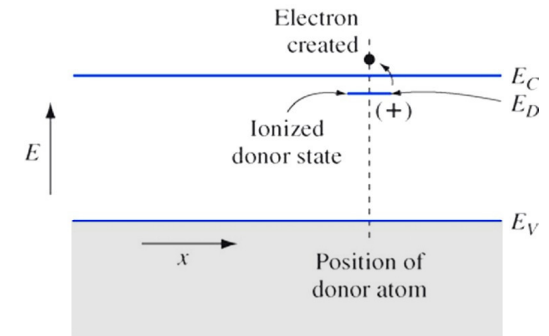
- ◆ Donor (or, Donor-like) level:

if filled by an electron → electrically neutral (電中性)
if empty → positively charged (帶+電)



- ◆ **Activation energy** (活化能量), or **ionization energy** (離子化能量)

$$E_d \equiv E_C - E_D$$

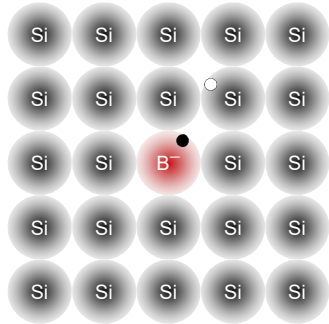


Extrinsic semiconductor: when an **acceptor** atom is present

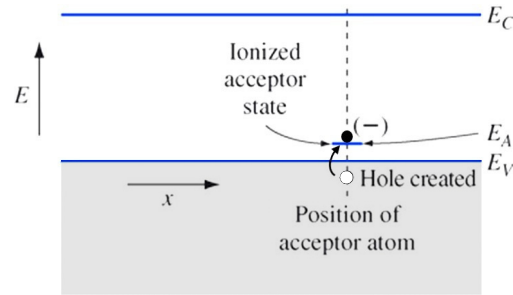


◆ Acceptor (or, Acceptor-like) level:

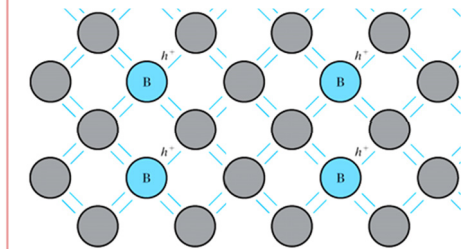
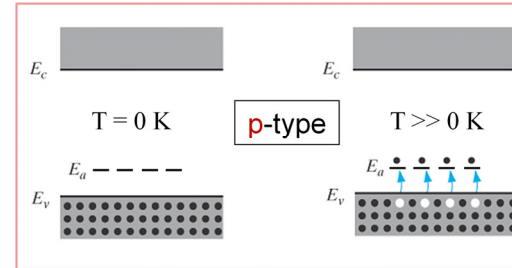
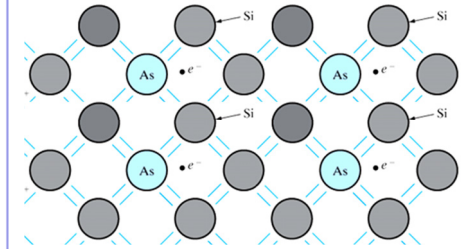
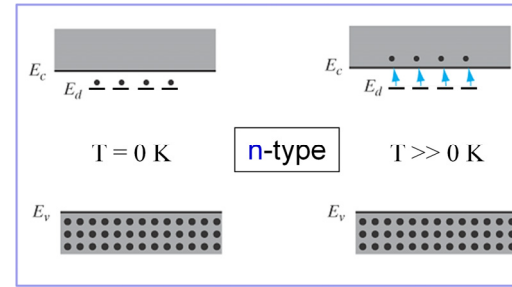
if filled by an electron → negatively charged (帶一電)
 if empty → electrically neutral (電中性)



◆ **Activation energy** (活化能量), or **ionization energy** (離子化能量)
 $E_a \equiv E_A - E_V$



When a large number of dopant atoms are present



Carrier Concentration at Thermal Equilibrium (熱平衡)



- ◆ Carrier concentration (載子濃度):
 Electron concentration (電子濃度), n
 Hole concentration (電洞濃度), p

◆ For a practical semiconductor:

$$n = N_C \cdot F_{1/2} \left(\frac{E_F - E_C}{kT} \right) \quad \& \quad p = N_V \cdot F_{1/2} \left(-\frac{E_F - E_V}{kT} \right)$$

where N_C is effective density of states in the conduction band (導電帶有效能態密度),
 N_V is effective density of states in the valence band (價電帶有效能態密度),
 $F_{1/2}(x)$ is Fermi integral (費米積分) with $n = 1/2$,

$$\left[\begin{array}{l} \text{Fermi integral: } F_n(x) = \frac{2}{\sqrt{\pi}} \int_0^{\infty} \frac{y^n}{1 + \exp(y-x)} dy \quad (\text{with } n > -1) \\ \text{In the case of calculating carrier concentration, we have } n = 1/2. \end{array} \right]$$

E_F is Fermi level (費米準位),
 k is Boltzmann's constant (波茲曼常數) ($= 8.62 \times 10^{-5} \text{ eV/K} = 1.38 \times 10^{-23} \text{ J/K}$),
 T is absolute temperature (絕對溫度).

$$n = N_C \cdot F_{1/2} \left(\frac{E_F - E_C}{kT} \right) \quad \text{where } N_C \text{ is effective density of states in the conduction band.}$$

$$p = N_V \cdot F_{1/2} \left(-\frac{E_F - E_V}{kT} \right) \quad \text{where } N_V \text{ is effective density of states in the valence band.}$$

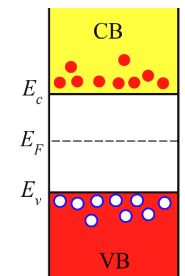
◆ If $E_C - 3kT > E_F > E_V + 3kT$, semiconductor is called **nondegenerate** (非簡併的).

$$n = N_C F_{1/2} \left(\frac{E_F - E_C}{kT} \right) \approx N_C \exp \left(\frac{E_F - E_C}{kT} \right)$$

$$p = N_V F_{1/2} \left(-\frac{E_F - E_V}{kT} \right) \approx N_V \exp \left(-\frac{E_F - E_V}{kT} \right)$$

$$\therefore np = N_V N_C \exp \left(-\frac{E_C - E_V}{kT} \right) = N_V N_C \exp \left(-\frac{E_g}{kT} \right)$$

$$\begin{array}{l} n = N_C e^{(E_F - E_C)/kT} \\ p = N_V e^{-(E_F - E_V)/kT} \\ n \cdot p = N_C N_V e^{-E_g/kT} \end{array}$$

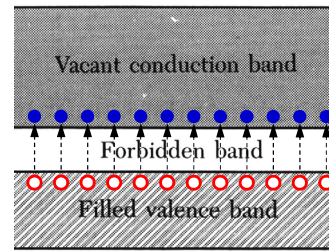
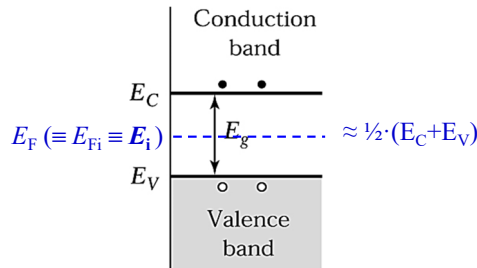


Intrinsic (i.e., undoped) semiconductors



- ◆ $n = p \equiv n_i$, where n_i is called **intrinsic carrier concentration** (本質載子濃度).
 $\Rightarrow n \cdot p = n_i \cdot n_i = n_i^2$
 But, an intrinsic semiconductor is nondegenerate. (See its E_F to be calculated below.)
 $\Rightarrow n \cdot p = N_C N_V e^{-E_g/kT}$
 \therefore The above two equations give $n_i = (N_C N_V)^{1/2} e^{-E_g/2kT}$.
- ◆ $n = p \Rightarrow N_C e^{(E_F - E_C)/kT} = N_V e^{-(E_F - E_V)/kT}$
 \therefore **Intrinsic Fermi level** (本質費米準位) is given by
 $E_F (\equiv E_{Fi} \equiv E_i) = \frac{1}{2} \cdot (E_C + E_V) + \frac{1}{2} \cdot kT \cdot \ln(N_V/N_C)$
 $\approx \frac{1}{2} \cdot (E_C + E_V)$

$$\begin{aligned} n &= N_C e^{(E_F - E_C)/kT} \\ p &= N_V e^{-(E_F - E_V)/kT} \\ n \cdot p &= N_C N_V e^{-E_g/kT} \\ &= n_i^2 \end{aligned}$$



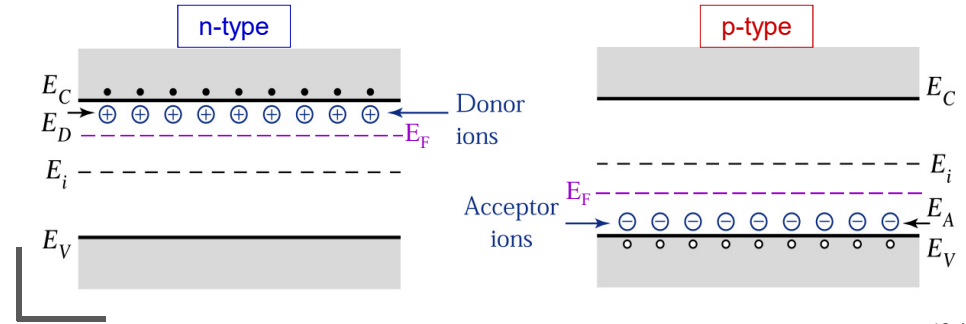
Extrinsic (i.e., doped) semiconductors



(N_D : donor concentration (施體濃度) & N_A : acceptor concentration (受體濃度))
 If N_D or N_A is **not very high**, giving the **semiconductor nondegenerate**, then we have:

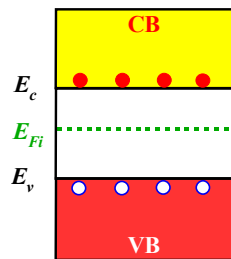
- ◆ **Complete ionization** (完全離子化): as shown in the figure below.
- ◆ If doped with N_D ($\gg n_i$) \Rightarrow **n-type**
 Majority carrier (多數載子) $n \approx N_D \Rightarrow E_F - E_C \approx kT \cdot \ln(N_D/N_C)$
 Minority carrier (少數載子) $p = N_V e^{-(E_F - E_V)/kT}$ or $p = n_i^2/n$
- ◆ If doped with N_A ($\gg n_i$) \Rightarrow **p-type**
 Majority carrier (多數載子) $p \approx N_A \Rightarrow E_F - E_V \approx kT \cdot \ln(N_V/N_A)$
 Minority carrier (少數載子) $n = N_C e^{(E_F - E_C)/kT}$ or $n = n_i^2/p$

$$\begin{aligned} n &= N_C e^{(E_F - E_C)/kT} \\ p &= N_V e^{-(E_F - E_V)/kT} \\ n \cdot p &= N_C N_V e^{-E_g/kT} \\ &= n_i^2 \end{aligned}$$



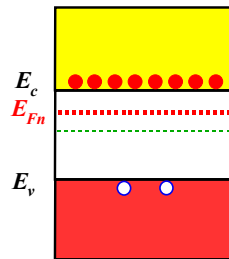
Intrinsic

$$\begin{aligned} n &= p \\ n \cdot p &= n_i^2 \end{aligned}$$



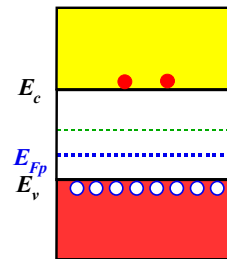
N-type

$$\begin{aligned} n &\gg p \\ n \cdot p &= n_i^2 \end{aligned}$$



P-type

$$\begin{aligned} n &\ll p \\ n \cdot p &= n_i^2 \end{aligned}$$



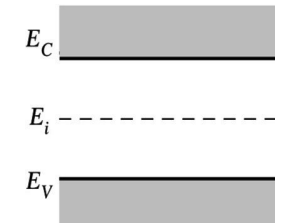
Summary of formulas (for both intrinsic & extrinsic cases)



When the **semiconductor is nondegenerate**:

- ◆ $n = N_C e^{(E_F - E_C)/kT} = N_C e^{(E_i - E_C)/kT} e^{(E_F - E_i)/kT} = n_i e^{(E_F - E_i)/kT}$
 $p = N_V e^{-(E_F - E_V)/kT} = N_V e^{-(E_i - E_V)/kT} e^{-(E_F - E_i)/kT} = n_i e^{-(E_F - E_i)/kT}$
 $n \cdot p = n_i^2 \leftarrow$ **Mass action law**
- ◆ $E_F - E_C = kT \cdot \ln(n/N_C)$
 $E_F - E_V = kT \cdot \ln(N_V/p)$
 or
 $E_F - E_i = kT \cdot \ln(n/n_i)$
 $E_F - E_i = kT \cdot \ln(n_i/p)$

$$\begin{aligned} n &= N_C e^{(E_F - E_C)/kT} \\ p &= N_V e^{-(E_F - E_V)/kT} \\ n \cdot p &= N_C N_V e^{-E_g/kT} \\ &= n_i^2 \end{aligned}$$



◆ For a doped semiconductor:

- If n type $\Rightarrow n \approx N_D$ & $p = n_i^2/n$
- If p type $\Rightarrow p \approx N_A$ & $n = n_i^2/p$

Symbols for undoping & doping concentration

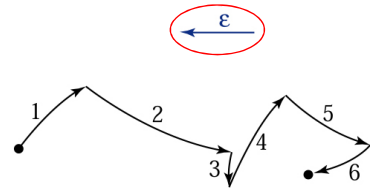
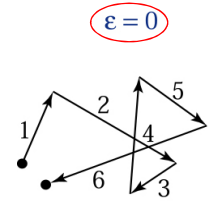
Examples: n^+ -Si, p^{++} -Si, n -GaAs, p^- -GaAs, i -GaAs

N_D : Donor concentration (施體濃度)
 N_A : Acceptor concentration (受體濃度)

N_D or N_A	Degenerate ?	N-type	P-type
$> 10^{19} \text{ cm}^{-3}$	Degenerate	n^{++}	p^{++}
$10^{18} \sim 10^{19} \text{ cm}^{-3}$	~ Degenerate	n^+	p^+
$10^{17} \sim 10^{18} \text{ cm}^{-3}$	Nondegenerate	n	p
$10^{15} \sim 10^{17}$	Nondegenerate	n^-	p^-
0 (Intrinsic)	Nondegenerate	i	i

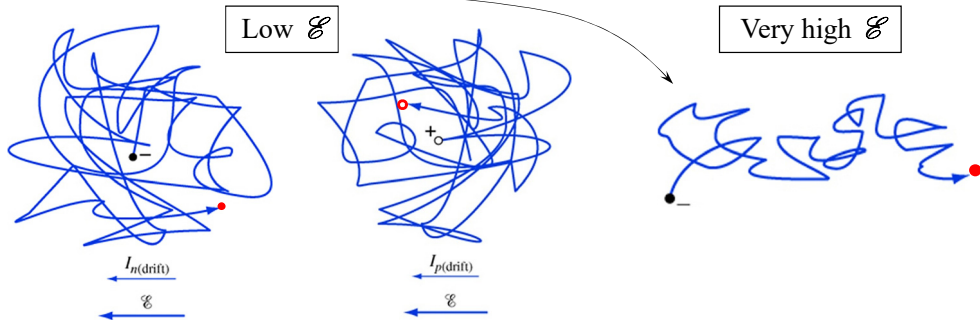
Drift (漂移)

- When $\mathcal{E} = 0$, carriers undergo random thermal motion only \Rightarrow Zero net displacement over a sufficiently long period of time.
- Average thermal velocity (平均熱速度), v_{th}
 $v_{th} \sim 10^7 \text{ cm/s @ 300 K}$
 (More precisely: $\frac{1}{2}m^*v_{th}^2 = 3 \cdot \frac{1}{2}kT$)
- Mean free path (平均自由路徑), l
 $l = v_{th} \cdot \tau_c$ ($\sim 1000 \text{ \AA @ 300 K}$)
- Mean free time (平均自由時間), or Collision time (碰撞時間), or Relaxation time (鬆弛時間)
 $\tau_c = l/v_{th}$ ($\sim 10^{-12} \text{ s} = 1 \text{ ps @ 300 K}$)
- When $\mathcal{E} \neq 0$, carriers have combined motion due to random thermal motion and an applied electric field. \Rightarrow Non-zero net displacement, so-called drift.



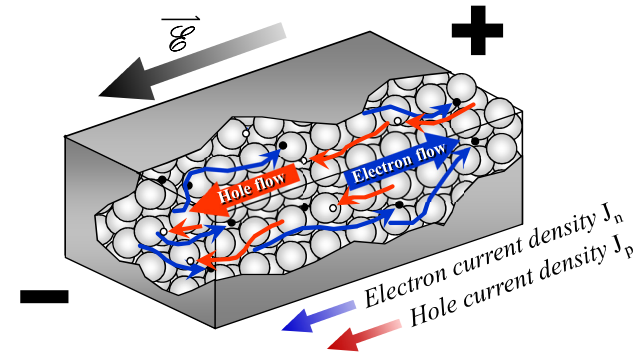
Drift Velocity (漂移速度) and Mobility (遷移率)

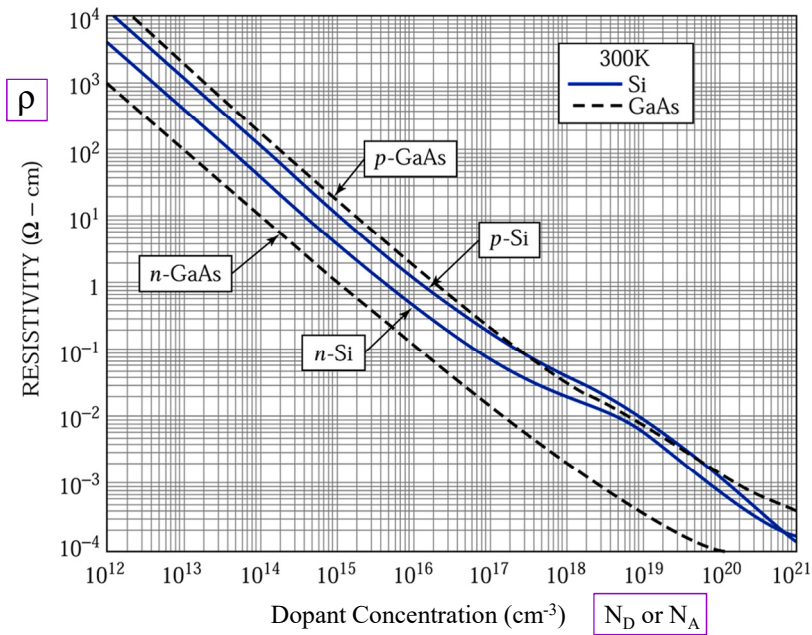
- Drift velocity (漂移速度) $\mathbf{v} = \mu \cdot \mathcal{E}$, where \mathcal{E} is the electric field.
- Mobility (遷移率) $\mu \equiv (q \cdot \tau_c) / m^*$, where m^* is the effective mass of the carriers.
 $(f \cdot t = m \cdot v \Rightarrow -q \mathcal{E} \cdot \tau_c = m \cdot v \Rightarrow v = -(q \tau_c / m) \mathcal{E})$
- In most of the semiconductors, electrons have higher mobility than holes (i.e., $\mu_n > \mu_p$) due to their smaller effective mass.
- Under low fields, drift velocity is much smaller than thermal speed (i.e., $v_d \ll v_{th}$).
- A trajectory such as this would take place only under very high fields.



Drift current (漂移電流)

- Drift current density (漂移電流密度): $\mathbf{J} = \mathbf{J}_n + \mathbf{J}_p$
 $\mathbf{J} = \mathbf{J}_n + \mathbf{J}_p = nq\mathbf{v}_n + pq\mathbf{v}_p = q(n\mu_n + p\mu_p) \mathcal{E} \equiv \sigma \cdot \mathcal{E}$
- Conductivity (導電係數 or 導電率): $\sigma = q(n\mu_n + p\mu_p)$ (in $\Omega^{-1} \cdot \text{cm}^{-1}$)
 - for n-type ($n \gg p$) $\Rightarrow \sigma \approx q \cdot n \cdot \mu_n$
 - for p-type ($p \gg n$) $\Rightarrow \sigma \approx q \cdot p \cdot \mu_p$
- Resistivity (電阻係數 or 電阻率): $\rho = 1/\sigma$ (in $\Omega \cdot \text{cm}$)





Resistivity versus doping concentration for Si and GaAs at 300 K.



Diffusion (擴散)

◆ If the carrier concentration is not uniformly distributed in space, the carriers can **diffuse** (擴散) even without experiencing an electric field.

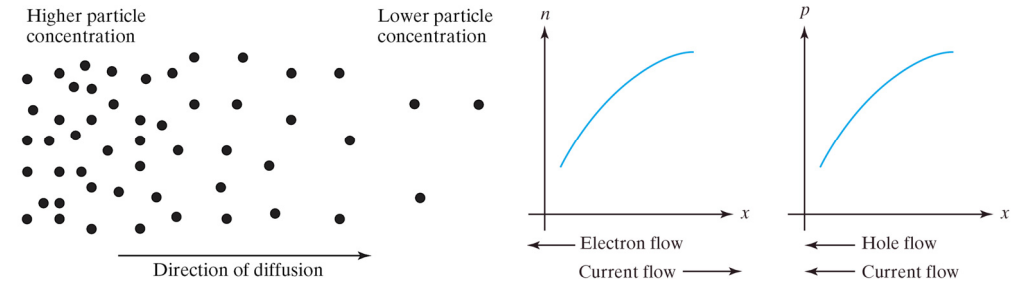
◆ **Diffusion current density** (擴散電流密度):

Electrons: $J_n = q \cdot D_n \cdot \frac{dn}{dx}$

$q = 1.6 \times 10^{-19} \text{ C}$

Holes: $J_p = -q \cdot D_p \cdot \frac{dp}{dx}$

where $D (= v_{th} \cdot l)$ is the **diffusion coefficient** (擴散係數) (or **diffusivity** (擴散率)).



Einstein Relation (愛因斯坦關係)

$$D = \left(\frac{kT}{q} \right) \cdot \mu$$

◆ $J_n = qD_n (\partial n / \partial x)$ is one-dimensional.

$$\therefore \frac{1}{2} m_n v_{th}^2 = \frac{1}{2} kT \Rightarrow v_{th}^2 = kT / m_n$$

◆ Use $\begin{cases} l = v_{th} \tau_c \\ \tau = \mu m^* / q \end{cases}$ ← from $\mu = q\tau / m^*$

$$\Rightarrow D_n = v_{th} l = v_{th} (v_{th} \tau_c) = v_{th}^2 \left(\frac{\mu_n m_n}{q} \right) = \left(\frac{kT}{m_n} \right) \left(\frac{\mu_n m_n}{q} \right)$$

$$\Rightarrow \left. \begin{aligned} D_n &= (kT/q) \mu_n && \text{for electrons} \\ D_p &= (kT/q) \mu_p && \text{for holes} \end{aligned} \right\} \text{Einstein relation}$$



High-Field Effects (高電場效應)

◆ At high electric fields, many effects may occur.

◆ Two important effects among them are:

- 1) **Drift velocity saturation** (漂移速度飽和)
- 2) **Impact ionization** (衝擊離子化)



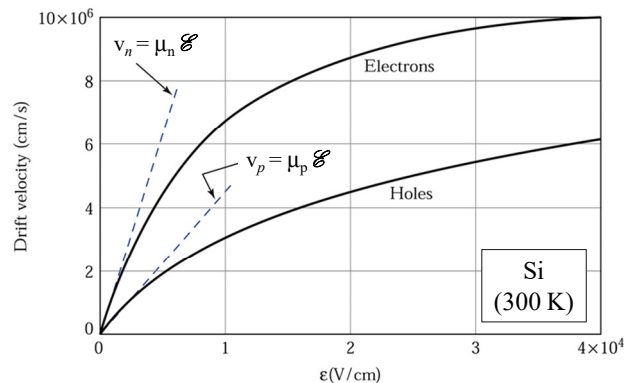


Drift Velocity Saturation (漂移速度飽和)

- Real drift velocity is represented by an empirical expression (經驗式)

$$v = \mu_0 \mathcal{E} / [1 + (\mu_0 \mathcal{E} / v_s)^\gamma]^{1/\gamma} \quad \leftarrow \text{for Si \& Ge}$$

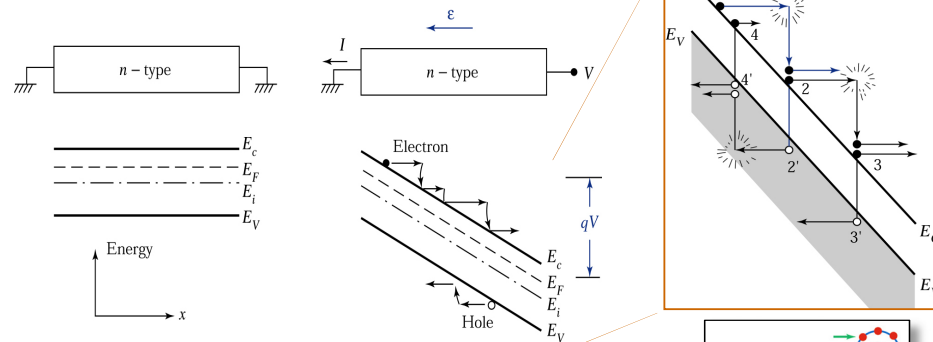
where μ_0 is the **low-field mobility** (低電場遷移率),
 v_s is the **saturation velocity** (飽和速度), and
 γ is a factor (2 for electrons; 1 for holes).



29, jswu

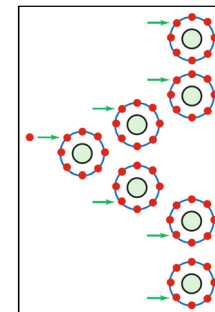


Impact Ionization (衝擊離子化)



- Impact ionization** (衝擊離子化), or **Avalanche multiplication** (累增放大)

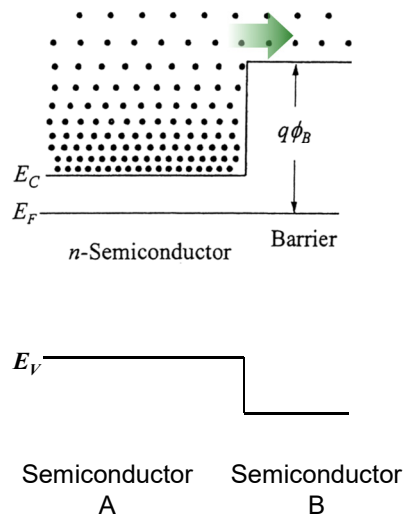
$$\begin{aligned} \frac{1}{2} m_1 v_s^2 &= E_g + 3 \cdot \frac{1}{2} m_1 v_f^2 \\ m_1 v_s &= 3 \cdot m_1 v_f \\ \Rightarrow \frac{1}{2} m_1 v_s^2 &= 1.5 E_g \\ \therefore \text{The required kinetic energy for the} \\ &\text{ionization process, } E_0 \approx 1.5 E_g. \end{aligned}$$



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Thermionic Emission (熱離發射)



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Tunneling (穿隧)

- In quantum mechanics:

Transmission coefficient (傳輸係數)

$$T = \left[1 + \frac{(qV_0 \sinh(\beta d))^2}{4E(qV_0 - E)} \right]^{-1}$$

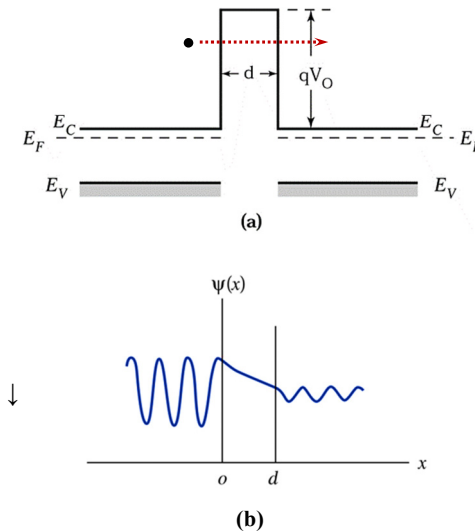
$$\text{where } \beta = \sqrt{2m_n(qV_0 - E)}/\hbar$$

- If $\beta \cdot d \gg 1$ (i.e., $V_0 \uparrow$, or $d \uparrow$, or $E \downarrow$)

$$\Rightarrow T \approx e^{-2\beta d} = e^{-2d \sqrt{2m_n(qV_0 - E)}/\hbar}$$

$\Rightarrow T$ is small

- If we want $T \uparrow$, we need $d \downarrow$, $qV_0 \downarrow$, $E \uparrow$, $m^* \downarrow$



(a) Band diagram of a semiconductor potential barrier with a distance d .

(b) Schematic representation of the wave function across the potential barrier.

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